

Fig. 4 | Extended optomechanical nanobeam lattice. **a**, Proposed experimental implementation of a topological nanobeam array. Nanobeams of different frequency are colour-coded, with grey rectangles depicting supports. Each nanobeam is interconnected to three nanobeams of different frequency via broadband nanophotonic cavities. **b**, This structure forms a honeycomb plaquette (surrounded by a dashed line), where each circle denotes an entire arm. By repetition, a honeycomb lattice with primitive lattice vectors is formed. The optical modulation phases to create a specific topological insulator phase, defined in the text, are indicated by the colours of the links. **c**, Simulated steady-state phononic amplitude under continuous driving, focused on the site indicated by an arrow. Chosen parameters are $\Gamma_1 = 2\pi = \Gamma_2 = 4$ kHz, $\Gamma_{\text{eff}} = 200$ kHz, phononic frequency disorder with standard deviation $2\pi = 20$ kHz and direct mechanical coupling of $t = 1$ kHz.

Non-trivial topological properties are then observed for rational ($= 1/3$ in the following), where the bandstructure for an infinite ribbon geometry shows q bulk bands split into q gaps, traversed by (chiral) counter-propagating edge modes (see top panel in Extended Data Fig. 2c). To exemplify phononic propagation along the boundary of a finite array, we calculate the steady state while driving an edge site (see Methods). A moderate frequency disorder of $\%$ is introduced, in addition to direct mechanical coupling of t_{ij} linking next-nearest neighbours²⁶. As shown by the calculated amplitudes in Fig. 4c, emerging phononic edge modes display unidirectional propagation within state-of-the-art fabrication tolerances (see Supplementary Section III and Extended Data Fig. 2c–e for more details). This demonstrates a feasible platform for nanoscale phononic topological insulators based on optically induced magnetic fields.

In conclusion, we have established a synthetic gauge field for nanomechanical transport using optically mediated couplings in an optomechanical system. We have introduced an experimental platform with large optomechanical interactions and bandwidths that unlocks many-mode implementations in the nanomechanical domain. The tunability of our system allows the investigation of

physics beyond the rotating wave approximation and in synthetic dimensions³¹. Out-of-plane optical control provides a counterpart to electromechanical methods^{32,33} with unique traits. It offers prospects for exploring the impact of thermal and quantum fluctuations on topologically protected phonon transport²¹ and the effects of mechanical or optomechanical nonlinearities²⁶. It is an advance towards exploiting topologically protected sound in the quantum acoustics regime and the search for exotic states such as those produced by non-Abelian gauge fields and analogues of the fractional quantum Hall effect³⁰ in the realm of nanomechanics.

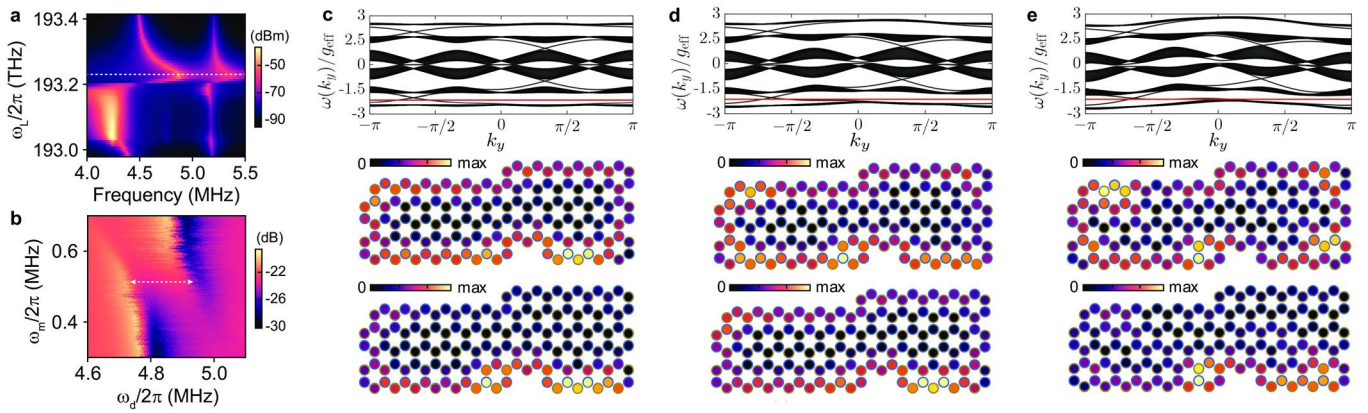
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Extended Data Fig. 2 | Robust phononic edge states. **a**, Optical spring shift measured at higher optical powers for a separate device shows large tunability of the mechanical modes. The dashed, white line shows the drive laser detuning used in **b**. The unusual features between 193.0 and 193.2 THz are due to a dynamically unstable regime of potential photothermoelastic origin. **b**, The modulated coupling strength, g_{eff} , is higher for higher optical powers. The largest modulated coupling strength measured in our experiments is marked by the dashed line and seen to be $\approx 2 \times 200$ kHz. Here the detection laser was absent and the drive laser response was directly demodulated using the lock-in amplifier leading to a Fano shaped feature for the driven response. **c,d,e**, (top panels) Band structure for a ribbon geometry (width $L = 20$) for increasing values of the direct mechanical coupling from left to right (0, 10, and 20 kHz, respectively), displaying the driving modulation frequency. The steady-state phononic amplitude in the absence of disorder is displayed in the middle row, while the result with a phononic frequency disorder with standard deviation $\sigma = 20$ kHz (averaged over 100 realizations) is shown at the bottom. For these plots, $\omega = \omega_m$.